

A high-order discontinuous Galerkin approach to the elasto-acoustic problem

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joint work with P. F. Antonietti and I. Mazzieri

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MODELLISTICA E CALCOLO SCIENTIFICO

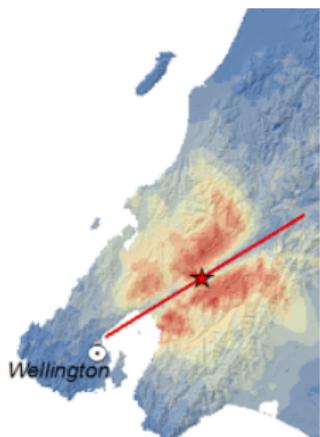
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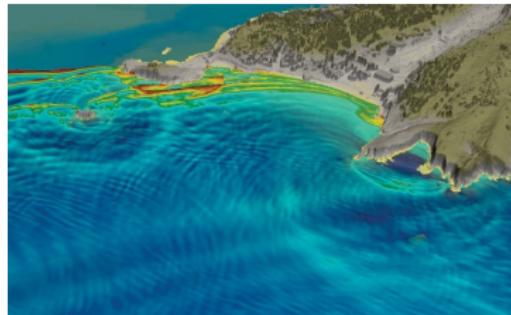
MODELING AND SCIENTIFIC COMPUTING

Motivations



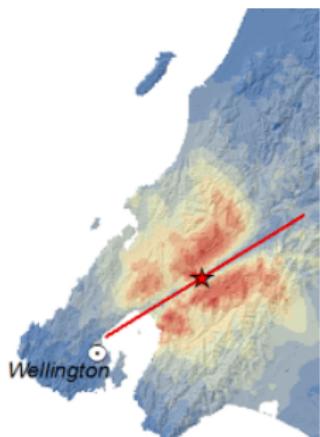
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- Simulation of **earthquake scenarios** near **coastal environments**
- Coupling of elastic and acoustic wave propagation



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Motivations

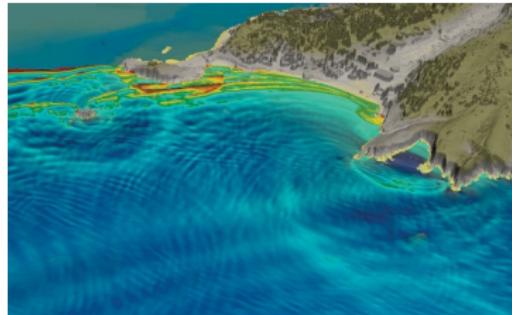


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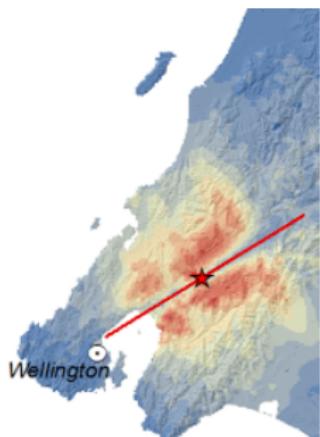
Requirements on the numerical scheme

- Flexibility
- Accuracy
- Efficiency

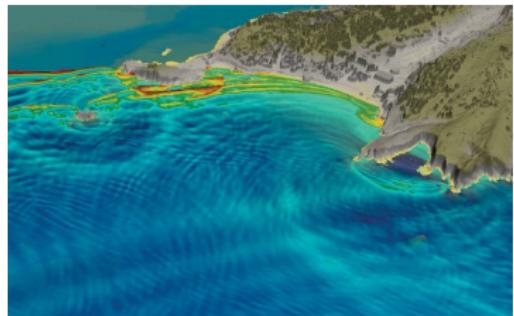


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Goal

- Numerical treatment based on **polyhedral meshes**
- The **dG method** supports **high-order polynomials** on such meshes

State of the art

Minimal bibliography

- [Komatitsch *et al.*, 2000]: Spectral Elements
- [Fischer and Gaul, 2005]: FEM–BEM coupling, Lagrange multipliers
- [Flemisch *et al.*, 2006]: classical FEM on two independent meshes
- [Brunner *et al.*, 2009]: FEM–BEM comparison
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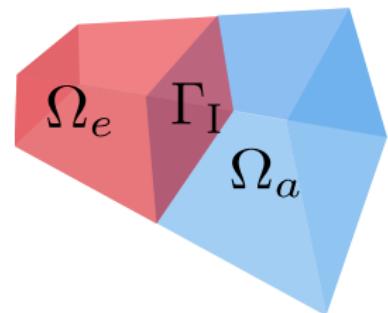
Our contribution

- **Well-posedness** of the coupled problem in the **continuous setting**
- **Detailed analysis** of a dG scheme on **general polytopic meshes**

The elasto-acoustic problem

Governing equations

$$\begin{cases} \rho_e \ddot{\mathbf{u}} + 2\rho_e \zeta \dot{\mathbf{u}} + \rho_e \zeta^2 \mathbf{u} - \operatorname{div} \boldsymbol{\sigma}(\mathbf{u}) = \mathbf{f}_e & \text{in } \Omega_e \times (0, T], \\ \boldsymbol{\sigma}(\mathbf{u}) - \mathbb{C}\boldsymbol{\varepsilon}(\mathbf{u}) = \mathbf{0} & \text{in } \Omega_e \times (0, T], \\ \mathbf{u} = \mathbf{0} & \text{on } \Gamma_{eD} \times (0, T], \\ \boldsymbol{\sigma}(\mathbf{u}) \mathbf{n}_e = -\rho_a \dot{\varphi} \mathbf{n}_e & \text{on } \Gamma_I \times (0, T], \\ \mathbf{u}(0) = \mathbf{u}_0, \quad \dot{\mathbf{u}}(0) = \mathbf{u}_1 & \text{in } \Omega_e, \\ c^{-2} \ddot{\varphi} - \Delta \varphi = f_a & \text{in } \Omega_a \times (0, T], \\ \varphi = 0 & \text{on } \Gamma_{aD} \times (0, T], \\ \partial \varphi / \partial \mathbf{n}_a = -\dot{\mathbf{u}} \cdot \mathbf{n}_a & \text{on } \Gamma_I \times (0, T], \\ \varphi(0) = \varphi_0, \quad \dot{\varphi}(0) = \varphi_1 & \text{in } \Omega_a \end{cases}$$



- The fluid exerts a pressure on the solid at the interface
- The normal component of the velocity is continuous at the interface

Well-posedness

Theorem (Existence and uniqueness)

Under suitable regularity hypotheses on initial data and source terms,
there is a **unique strong solution** s.t.

$$\begin{aligned}\boldsymbol{u} &\in C^2([0, T]; \boldsymbol{L}^2(\Omega_e)) \cap C^1([0, T]; \boldsymbol{H}_D^1(\Omega_e)) \cap C^0([0, T]; \boldsymbol{H}_{\mathbb{C}}^\Delta(\Omega_e) \cap \boldsymbol{H}_D^1(\Omega_e)), \\ \varphi &\in C^2([0, T]; L^2(\Omega_a)) \cap C^1([0, T]; H_D^1(\Omega_a)) \cap C^0([0, T]; H^\Delta(\Omega_a) \cap H_D^1(\Omega_a)),\end{aligned}$$

$$\begin{aligned}\boldsymbol{H}_{\mathbb{C}}^\Delta(\Omega_e) &:= \{\boldsymbol{v} \in \boldsymbol{L}^2(\Omega_e) : \operatorname{div} \mathbb{C}\boldsymbol{\epsilon}(\boldsymbol{v}) \in \boldsymbol{L}^2(\Omega_e)\}, \\ H^\Delta(\Omega_a) &:= \{v \in L^2(\Omega_a) : \Delta v \in L^2(\Omega_a)\}\end{aligned}$$

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Proof

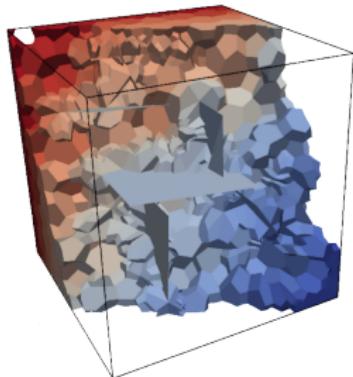
Apply **Hille–Yosida** upon rewriting the system as

$$\begin{aligned}\frac{d\mathcal{U}}{dt}(t) + A\mathcal{U}(t) &= \mathcal{F}(t), \quad t \in (0, T], \\ \mathcal{U}(0) &= \mathcal{U}_0\end{aligned}$$

Meshes and spaces

Mesh

- Nonconforming **polyhedral** mesh $\mathcal{T}_h = \mathcal{T}_h^e \cup \mathcal{T}_h^a$
- **Arbitrary** number of faces per element
- Possible presence of **degenerating faces**

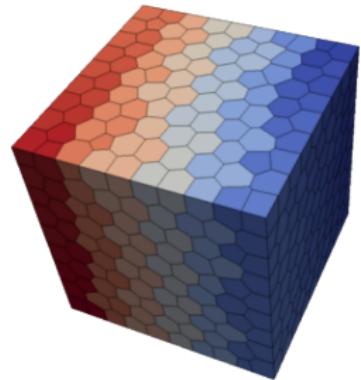


[Cangiani *et al.*, 2017], [Antonietti *et al.*, 2017]

Discrete spaces

$$V_h^e = \{\boldsymbol{v}_h \in \mathbf{L}^2(\Omega_e) : \boldsymbol{v}_{h|\kappa} \in [\mathcal{P}_{p_{e,\kappa}}(\kappa)]^d \ \forall \kappa \in \mathcal{T}_h^e\},$$

$$V_h^a = \{\psi_h \in L^2(\Omega_a) : \psi_{h|\kappa} \in \mathcal{P}_{p_{a,\kappa}}(\kappa) \ \forall \kappa \in \mathcal{T}_h^a\}$$



Semi-discrete problem

Find $(\mathbf{u}_h, \varphi_h) \in C^2([0, T]; \mathbf{V}_h^e) \times C^2([0, T]; V_h^a)$ s.t.,
for all $(\mathbf{v}_h, \psi_h) \in \mathbf{V}_h^e \times V_h^a$,

$$\begin{aligned} & (\rho_e \ddot{\mathbf{u}}_h(t), \mathbf{v}_h)_{\Omega_e} + (c^{-2} \rho_a \ddot{\varphi}_h(t), \psi_h)_{\Omega_a} \\ & + (2\rho_e \zeta \dot{\mathbf{u}}_h(t), \mathbf{v}_h)_{\Omega_e} + (\rho_e \zeta^2 \mathbf{u}_h(t), \mathbf{v}_h)_{\Omega_e} \\ & + \mathcal{A}_h^e(\mathbf{u}_h(t), \mathbf{v}_h) + \mathcal{A}_h^a(\varphi_h(t), \psi_h) \\ & + \mathcal{I}_h^e(\dot{\varphi}_h(t), \mathbf{v}_h) + \mathcal{I}_h^a(\dot{\mathbf{u}}_h(t), \psi_h) \\ & = (\mathbf{f}_e(t), \mathbf{v}_h)_{\Omega_e} + (f_a(t), \psi_h)_{\Omega_a} \end{aligned}$$

Semi-discrete problem

Find $(\mathbf{u}_h, \varphi_h) \in C^2([0, T]; \mathbf{V}_h^e) \times C^2([0, T]; V_h^a)$ s.t.,
for all $(\mathbf{v}_h, \psi_h) \in \mathbf{V}_h^e \times V_h^a$,

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Stability

For sufficiently large stabilization parameters,

$$\|(\mathbf{u}_h(t), \varphi_h(t))\|_{\mathcal{E}} \lesssim \|(\mathbf{u}_h(0), \varphi_h(0))\|_{\mathcal{E}} + \int_0^t (\|\mathbf{f}_e(\tau)\|_{\Omega_e} + \|f_a(\tau)\|_{\Omega_a}) d\tau$$

Error estimate

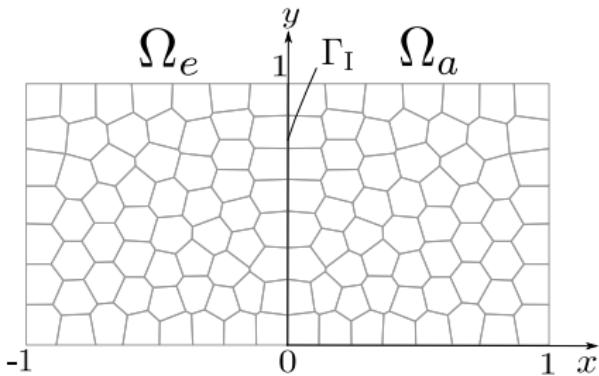
Energy error estimate

If $(\mathbf{u}, \varphi) \in C^2([0, T]; \mathbf{H}^m(\Omega_e)) \times C^2([0, T]; H^n(\Omega_a))$, with $m \geq p_e + 1$ and $n \geq p_a + 1$,
and if stabilization parameters are sufficiently large,

$$\begin{aligned} \sup_{t \in [0, T]} \|(\mathbf{e}_e(t), e_a(t))\|_{\mathcal{E}}^2 &\lesssim \frac{h^{2p_e}}{p_e^{2m-3}} \left(\sup_{t \in [0, T]} \left(\|\dot{\mathbf{u}}(t)\|_{m, \Omega_e}^2 + \|\mathbf{u}(t)\|_{m, \Omega_e}^2 \right) \right. \\ &\quad \left. + \int_0^T \left(\|\ddot{\mathbf{u}}(t)\|_{m, \Omega_e}^2 + \|\dot{\mathbf{u}}(t)\|_{m, \Omega_e}^2 + \|\mathbf{u}(t)\|_{m, \Omega_e}^2 \right) dt \right) \\ &\quad + \frac{h^{2p_a}}{p_a^{2n-3}} \left(\sup_{t \in [0, T]} \left(\|\dot{\varphi}(t)\|_{n, \Omega_a}^2 + \|\varphi(t)\|_{n, \Omega_a}^2 \right) \right. \\ &\quad \left. + \int_0^T \left(\|\ddot{\varphi}(t)\|_{n, \Omega_a}^2 + \|\dot{\varphi}(t)\|_{n, \Omega_a}^2 + \|\varphi(t)\|_{n, \Omega_a}^2 \right) dt \right) \end{aligned}$$

Optimal in h , suboptimal in p by a factor $1/2$

Numerical example I



Test case 1

We solve the elasto-acoustic problem on $\Omega_e \cup \Omega_a$, for $T = 1$, $\Delta t = 10^{-4}$, for a homogeneous isotropic elastic material, such that

$$\begin{aligned}\mathbf{u}(x, y; t) &= x^2 \cos(\sqrt{2}\pi t) \cos\left(\frac{\pi}{2}x\right) \sin(\pi y) \hat{\mathbf{u}}, & \hat{\mathbf{u}} &= (1, 1) \\ \varphi(x, y; t) &= x^2 \sin(\sqrt{2}\pi t) \sin(\pi x) \sin(\pi y),\end{aligned}$$

Numerical example I

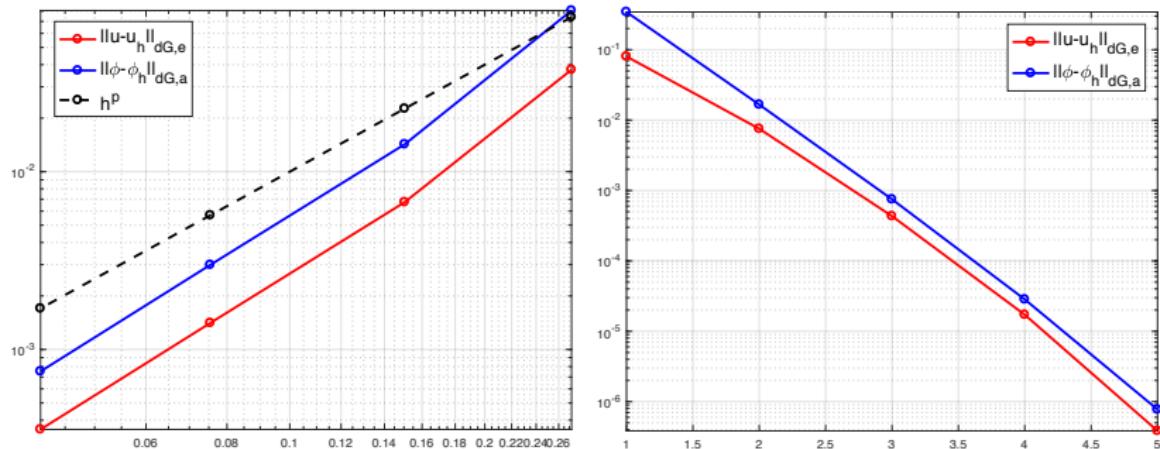
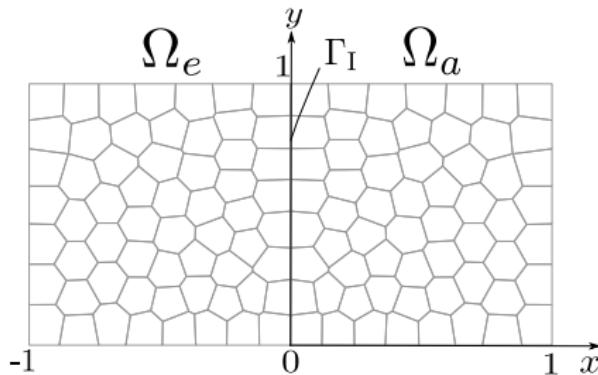


Figure: $\|u - u_h\|_{dG,e}$ and $\|\varphi - \varphi_h\|_{dG,a}$ vs. h (left) and p (right) at $T = 1$

Numerical example II



Test case 2 [Mönköla, 2016]

We solve the elasto-acoustic problem on $\Omega_e \cup \Omega_a$, for $T = 0.8$, $\Delta t = 10^{-4}$, for a homogeneous isotropic elastic material, such that

$$\begin{aligned}\mathbf{u}(x, y; t) &= \left(\cos\left(\frac{4\pi x}{c_p}\right), \cos\left(\frac{4\pi x}{c_s}\right) \right) \cos(4\pi t), \\ \varphi(x, y; t) &= \sin(4\pi x) \sin(4\pi t),\end{aligned} \quad c_p = \sqrt{\frac{\lambda + 2\mu}{\rho_e}}, \quad c_s = \sqrt{\frac{\mu}{\rho_e}}$$

Numerical example II

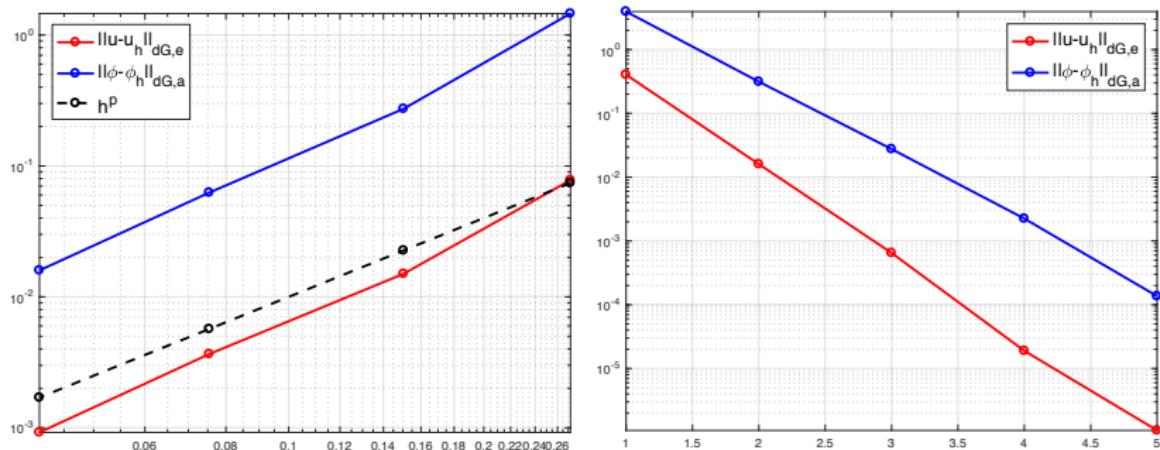
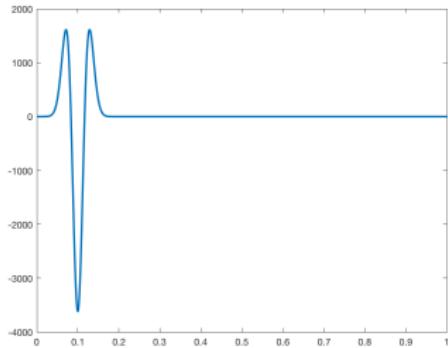
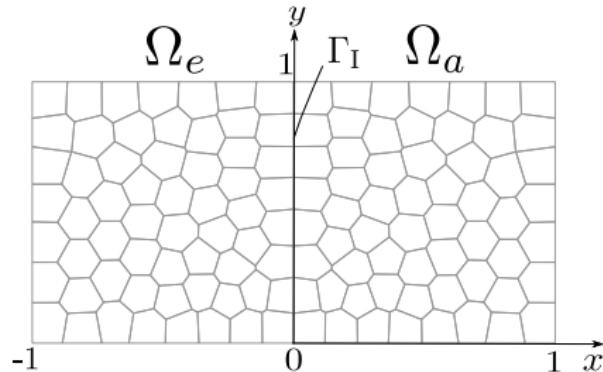


Figure: $\|u - u_h\|_{dG,e}$ and $\|\varphi - \varphi_h\|_{dG,a}$ vs. h (left) and p (right) at $T = 0.8$

Physical example



$$t \mapsto -2\pi a \left(1 - 2\pi a(t - t_0)^2\right) e^{-\pi a(t - t_0)^2}$$

Physical example

We simulate a seismic source in the acoustic domain by a point-wise source:

$$f_a(\mathbf{x}, t) = -2\pi a \left(1 - 2\pi a(t - t_0)^2\right) e^{-\pi a(t - t_0)^2} \delta(\mathbf{x} - \mathbf{x}_0), \quad \mathbf{x}_0 \in \Omega_a, \quad t_0 \in (0, T],$$

$$\mathbf{x}_0 = (0.2, 0.5), \quad t_0 = 0.1$$

Physical example

$t \mapsto \|\mathbf{u}(x, y; t)\|_2$ and $t \mapsto |\varphi(x, y; t)|$

Conclusions & perspectives

Conclusions

- We proved that the elasto-acoustic problem is well-posed in the continuous setting
- We proved and numerically validated hp -convergence results for a discontinuous Galerkin method on polyhedral meshes
- We used the method to simulate an example of physical interest

Conclusions & perspectives

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- We proved that the elasto-acoustic problem is well-posed in the continuous setting
- We proved and numerically validated hp -convergence results for a discontinuous Galerkin method on polyhedral meshes
- We used the method to simulate an example of physical interest

Perspectives

- Carrying out numerical computations in a 3D setting, using the code SPEED (<http://speed.mox.polimi.it/>)
- Considering the case of totally absorbing boundary conditions
- Deducing error estimates for the fully discrete problem
- Consider the more general case of a viscoelastic material response:

$$\boldsymbol{\sigma}(\boldsymbol{u}(\boldsymbol{x}, t); t) = \mathbb{C}(\boldsymbol{x}, 0)\boldsymbol{\varepsilon}(\boldsymbol{u}(\boldsymbol{x}, t)) - \int_0^t \frac{\partial \mathbb{C}}{\partial s}(\boldsymbol{x}, t-s)\boldsymbol{\varepsilon}(\boldsymbol{u}(\boldsymbol{x}, s)) \, ds$$

References I

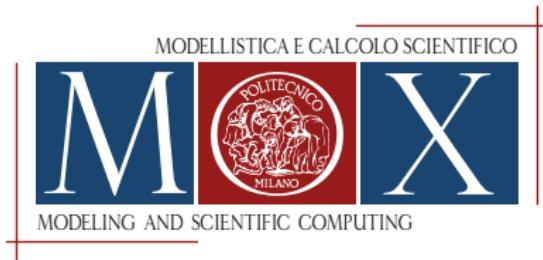
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Merci de votre attention



Application of Hille–Yosida

Let $\mathbf{w} = \dot{\mathbf{u}}$, $\phi = \dot{\varphi}$, and $\mathcal{U} = (\mathbf{u}, \mathbf{w}, \varphi, \phi)$. We introduce

$$\mathbb{H} = \mathbf{H}_D^1(\Omega_e) \times \mathbf{L}^2(\Omega_e) \times H_D^1(\Omega_a) \times L^2(\Omega_a),$$

with scalar product

$$\begin{aligned} (\mathcal{U}_1, \mathcal{U}_2)_{\mathbb{H}} &= (\rho_e \zeta^2 \mathbf{u}_1, \mathbf{u}_2)_{\Omega_e} + (\mathbb{C} \boldsymbol{\varepsilon}(\mathbf{u}_1), \boldsymbol{\varepsilon}(\mathbf{u}_2))_{\Omega_e} \\ &\quad + (\rho_e \mathbf{w}_1, \mathbf{w}_2)_{\Omega_e} + (\rho_a \nabla \varphi_1, \nabla \varphi_2)_{\Omega_a} + (c^{-2} \rho_a \phi_1, \phi_2)_{\Omega_a}. \end{aligned}$$

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Then, we define the operator $A: D(A) \subset \mathbb{H} \rightarrow \mathbb{H}$ by

$$A\mathcal{U} = (-\mathbf{w}, 2\zeta\mathbf{w} + \zeta^2 \mathbf{u} - \rho_e^{-1} \operatorname{div} \mathbb{C}\boldsymbol{\varepsilon}(\mathbf{u}), -\phi, -c^2 \Delta \varphi) \quad \forall \mathcal{U} \in D(A),$$

$$\begin{aligned} D(A) = \Big\{ \mathcal{U} \in \mathbb{H} : & \mathbf{u} \in \mathbf{H}_{\mathbb{C}}^{\triangle}(\Omega_e), \mathbf{w} \in \mathbf{H}_D^1(\Omega_e), \varphi \in H^{\triangle}(\Omega_a), \phi \in H_D^1(\Omega_a); \\ & (\mathbb{C}\boldsymbol{\varepsilon}(\mathbf{u}) + \rho_a \phi \mathbf{I}) \mathbf{n}_e = \mathbf{0} \text{ on } \Gamma_I, \quad (\nabla \varphi + \mathbf{w}) \cdot \mathbf{n}_a = 0 \text{ on } \Gamma_I \Big\}. \end{aligned}$$

Application of Hille–Yosida

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with scalar product

$$\begin{aligned} (\mathcal{U}_1, \mathcal{U}_2)_{\mathbb{H}} &= (\rho_e \zeta^2 \mathbf{u}_1, \mathbf{u}_2)_{\Omega_e} + (\mathbb{C}\boldsymbol{\varepsilon}(\mathbf{u}_1), \boldsymbol{\varepsilon}(\mathbf{u}_2))_{\Omega_e} \\ &\quad + (\rho_e \mathbf{w}_1, \mathbf{w}_2)_{\Omega_e} + (\rho_a \nabla \varphi_1, \nabla \varphi_2)_{\Omega_a} + (c^{-2} \rho_a \phi_1, \phi_2)_{\Omega_a}. \end{aligned}$$

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Finally, let $\mathcal{F} = (\mathbf{0}, \rho_e^{-1} \mathbf{f}_e, 0, c^2 f_a)$.

For $\mathcal{F} \in C^1([0, T]; \mathbb{H})$ and $\mathcal{U}_0 \in D(A)$,
find $\mathcal{U} \in C^1([0, T]; \mathbb{H}) \cap C^0([0, T]; D(A))$ s.t.

$$\frac{d\mathcal{U}}{dt}(t) + A\mathcal{U}(t) = \mathcal{F}(t), \quad t \in (0, T],$$

$$\mathcal{U}(0) = \mathcal{U}_0.$$